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KRISTIAN HERTZ

SIMPLE TEMPERATURE CALCULATIONS OF FIRE EXPOSED CONCRETE CONTRUCTIONS

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Joseph Fourier

PREFACE

Joseph Fourier wrote the comprehensive treatise:
"Théorie du Mouvement de la Chaleur dans les Corps Solides" [3] only a few years after his return from Egypt, where he had participated in the campaign of Napoleon Bonaparte in 1799.

The work has been a valuable part of the natural philosophy ever since and a direct source of inspiration for scientists as J.C. Maxwell and Lord Kelvin.

Especially the following theorem has become of outstanding importance to fire technology: "Suppose the different points of a homogeneous solid of any form whatever, to have received initial temperatures which vary successively by the effect of the mutual action of the molecules, and suppose the equation v = f(x,y,z,t) to represent the successive states of the solid, it may now be shown that v a function of four variables necessarily satisfies the equation

$$\frac{dv}{dt} = \frac{K}{CD} \left(\frac{d^2v}{dx^2} + \frac{d^2v}{dy^2} + \frac{d^2v}{dz^2} \right) .$$

This report deals with solutions to the equation.

Copenhagen, June 1981 Kristian Hertz M.Sc. Ph.D. Struct. Eng.

SUMMARY

The problem of determining the temperature d: bution in fire exposed concrete constructions analysed.

Based on the knowledge of the thermal propertion concrete the accuracy of a calculational process estimated leading to the conclusion, that application of simple approximative solutions well justified.

Some simple exact solutions to Fouriers equal presented, and a new procedure is developed obing the temperature distribution in a rectar concrete specimen exposed to a realistic fire standard fire course extended with a decay per

In four appendices the method is formulated : pocket calculator program and a fast EDP subsand examples of calculated distributions are with results of more complicated calculations measurements from fire tests.

ACKNOWLEDGEMENTS

I would like to express my gratitude to senic turer T. Jakobsen for reading the manuscript staff of the Institute of Building Design for and printing the report.

SYMBOLS

a _l	constant
a	thermal diffusivity
С	heat capacity as used by Fourier
Ci	constants
Ci	half period
c _p	specific capacity of heat
D	density as used by Fourier
D,	thermal amplitude
E'	constant temperature
f	functions
i	enthalpy
K	thermal conductivity
L	constant
m	constant
r	radius
T	temperature
t	time
v	angular velocity
v	temperature as used by Fourier
x	coordinate
У	coordinate
z	coordinate
$\xi_{\mathbf{T}\mathbf{x}}$	temperature reduction in the depth x
ρ	density

INTRODUCTION

A precondition for an analytical determination of the fire resistance of a construction is the abilof using a calculational approach for estimating temperature distribution through the loadbearing pof it.

While the laws of statics remain unchanged the properties of the materials are sensitive to the temperature development.

The temperature is the key parameter, and its varriation as a function of time and place is especia important for the understanding of the function of fire exposed concrete structures, because large thermal differences normally occur, giving rise to internal stresses and a lack of simultaneity which is often ignored without any reason.

In an age of electronic data processing the standa answer to complicated questions is calculational power. Nevertheless simple procedures are still ju fied in order to achieve reasonable solutions quic and at small efforts.

The need for procedures of this kind for calculati of temperatures in fire exposed constructions is i creasing because still more constructions has to b designed for resistance to fire.

This is the case especially if the thermal propert which are used in the calculations are so poorly d termined, that a greater accuracy of the procedure is meaningless.

ON THE THERMAL CONDUCTIVITY OF CONCRETE

The sort of aggregate used is highly decisive for the value of the thermal conductivity and its variation with the temperature.

The water-cement ratio, and dependent on this, the porosity is also of major importance.

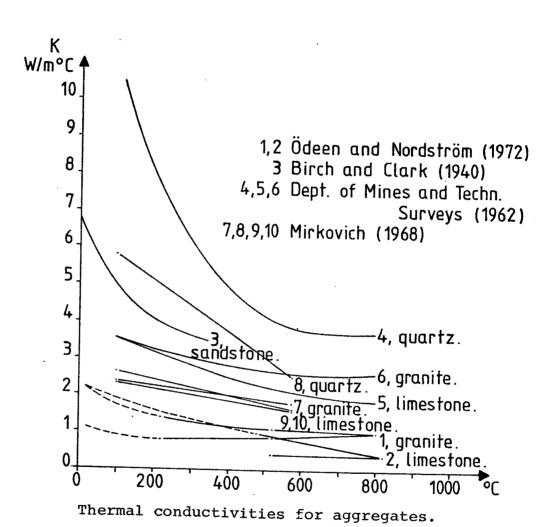
It is therefore obvious that if the thermal conductivity is not measured for the actual concete the assessment of this value can give rise to considerable deviations when calculating a temperature distribution.

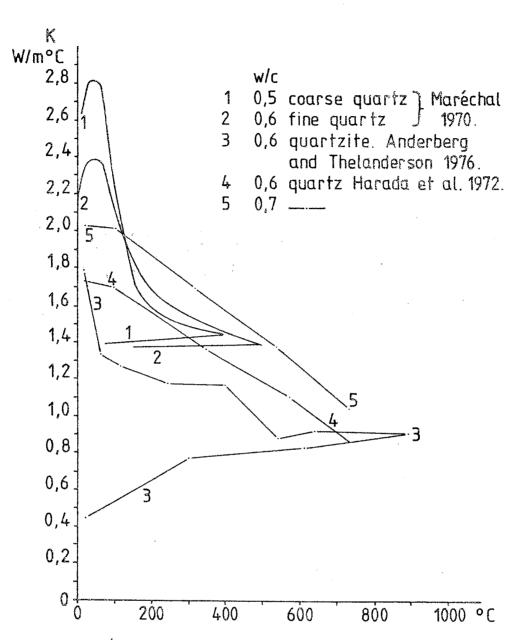
If the value is measured as a function of the temperature, the conditions under which such a measurement is performed are much different from the conditions being found in a construction exposed to fire. The concrete is in the former case mostly dry and the duration of the heating is often several days.

Thus, no matter how careful the conductivity is determined and no matter how precise the tempeature calculation has been executed there will always be an uncertainty of say 40- to 50° C on the temperature distribution in a fire exposed cross section solely arising from uncertainty on the determination of the thermal conductivity.

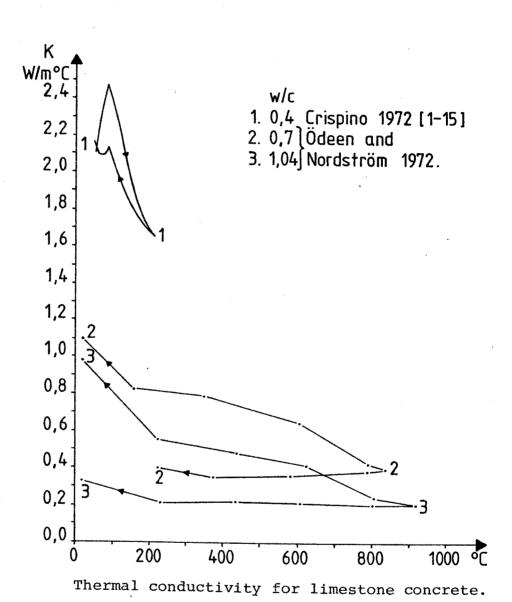
Some examples of the variation of the thermal conductivity by the temperature are shown on the following pages.

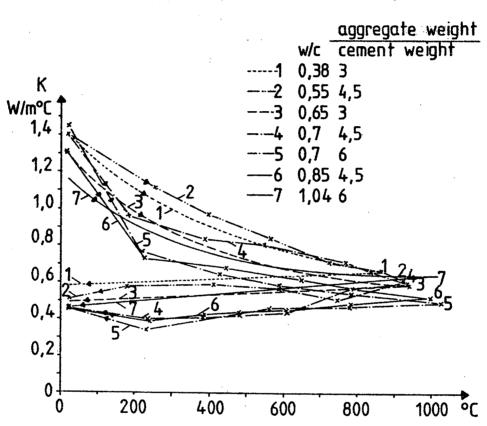
To be noticed is the variation caused by the differences in aggregate and the influence of the heating rate.



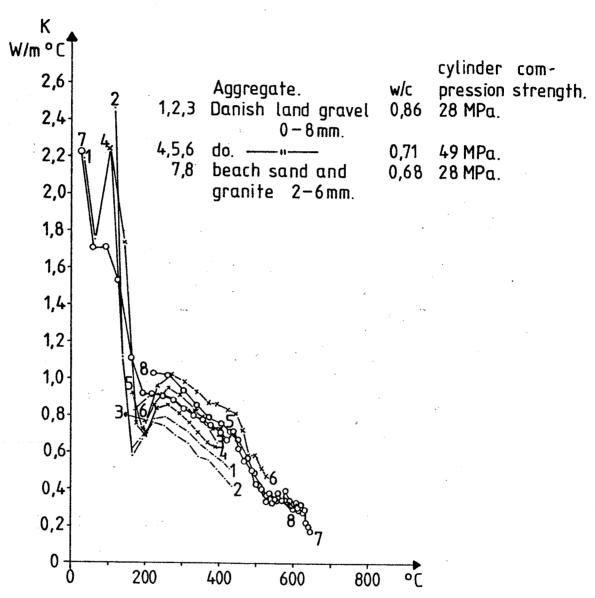


Thermal conductivity for concrete with quartz.





Thermal conductivity for granite concrete. Ödeen and Nordström [18].



Thermal conductivity during a standard fire. Østergaard 1972 [21].

ON THE SPECIFIC CAPACITY OF HEAT

Unlike the thermal conductivity the specific heat capacity of dry concrete seems to vary only a little with differences in the composition and with the temperature level.

Many authors like Lie [7] and Pettersson and Ödeen [12] propose to use a fixed value of approximately $1.0\,\mathrm{kJ/kg^OC}$, and the results of the thorough work of Ödeen and Nordström [18] seems to confirm the reasonableness of such an assumption, as the actual value raises only 5- to 10 pct, when the temperature raises from $200\,\mathrm{^OC}$ to $600\,\mathrm{^OC}$.

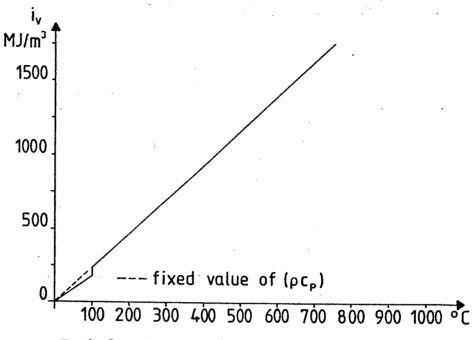
During the experimental investigations the concrete has been examined while cooling and not while heating, which would be more interesting for these applications

Although all processes involved for the dry concrete can not be provided to be reversible at the same temperature levels many authors considers the deviations caused by this to be negligible.

On the other hand one must of course correct for the influence of moisture on the heat demand of the concrete supposed to be heated, especially the local increase of the heat capacity at 100- to 200°C due to evaporation of the free water.

For very precise calculations the moisture can be handled separately taking into account the properties of heat and flow.

l kg water uses 2.6 MJ while heated from 20- to 100° C and then finally evaporated.



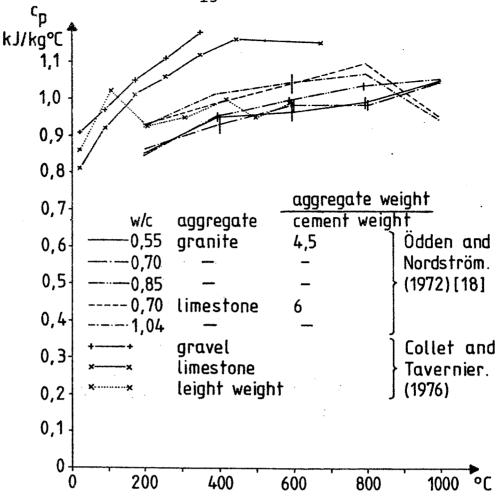
Enthalpy in principle.

The moisture content of structural concrete is ofter about 1- to 3 pct. by weight. That is about 1 pct. for protected structures and 2- to 3 pct. for structures more likely to be exposed to moisture. (See fc example Neville [10] p. 429). A realistic value thus is about 1.5 weight-pct. free water.

While heated to 200° C this moisture uses about $1.5x^{2}$ ~ 40 kJ per kg concrete, and related to the total heat demand of the concrete, which is about 200 kJ per kg concrete, this is approximately 20 pct.

An examination of curves showing the development of the specific heat capacity for actual examples of dr concrete indicates that the value is just about 20 r less than 1.0 kJ/kg^OC within the first 200^OC.

This means that the value 1.0 kJ/kg C seems to de-



Specific capacities of heat.

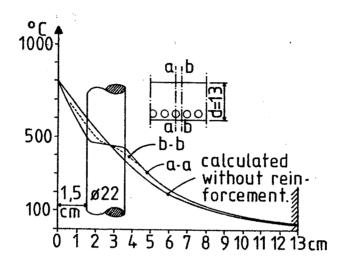
scribe the heat capacity for the total system almost exatly at all temperatures when used in relation to the density of the dry concrete.

Above 200°C the density of the dry concrete decreases slightly.

Using the results from Zoldners [17] the density has decreased 5 pct. at 600°C and thus about the same amount as the slight increase in the specific heat capacity above 200°C. So the product of these two values must be expected to remain almost constant.

This phenomenon facilitates the simple calculations.

CALCULATED TEMPERATURE DISTRIBUTIONS



Temperature distribution in a 13 cm slab after 1 h standard fire.

A result of major importance to this subject was published by Ehm in 1967, where he showed that if the temperature distribution in a reinforced concret section is calculated as if the section consists of plain concrete, the temperatures at the positions of the centres of the reinforcement bars will be the same as the temperatures of the bars in the corresponding reinforced cross section.

Becker et al. [1] showed that this is a reasonable procedure up to an area of reinforcement of 4% of the total cross sectional area.

In most cases the problem therefore is reduced to the calculation of the temperature distribution in a pla concrete section.

If the effect of evaporation and moisture flow is neglected the heat transport is ruled by the differential equation - Fourier's law:

$$\frac{\partial i_{v}}{\partial t} = \nabla (K \nabla T)$$

where i_v is the enthalpy (J/m^3) , t is the time (s), T is the temperature (for example ^{O}C) and K is the thermal conductivity $(W/m^{O}C)$.

The corresponding one-dimensional expression is then with fixed conductivity K, specific capacity of heat c_p (J/kg^OC) and density ρ (kg/m³):

$$\frac{\partial^2 \mathbf{T}}{\partial \mathbf{x}^2} = \frac{\rho \mathbf{c}_{\mathbf{p}}}{K} \frac{\partial \mathbf{T}}{\partial \mathbf{t}}$$

where x is a simple Euclidean co-ordinate.

Dividing the cross-section into slices of thickness Δx the corresponding difference expression can be solved stepwise graphically or by EDP.

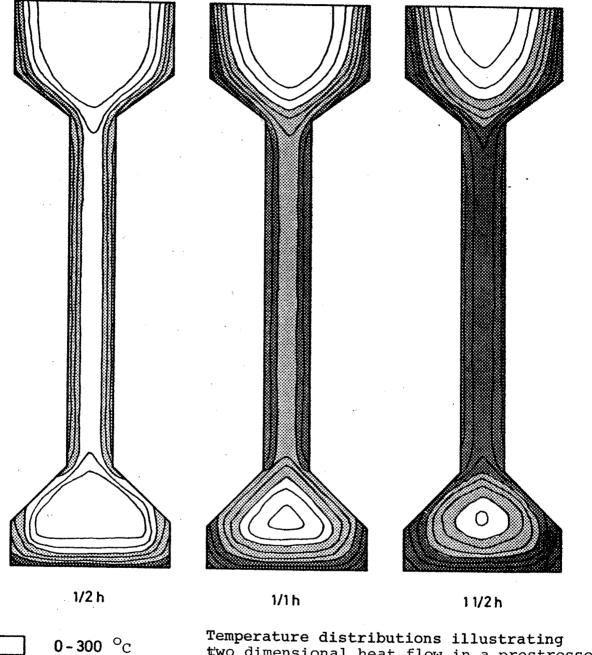
Results of the latter kind are published by for example Ödeen [20], Lie [7] and Maes et al. [8].

Also for two-dimensional heat flow difference expressions are developed and stepwise solutions have been found.

Such solutions are available in for example Odeen [19], Weiss [13] and Pettersson and Odeen [12].

Also the finite element method is applicable to the problem using an appropriate principle of variation.

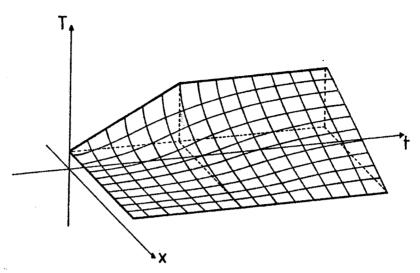
See for example Zienkiewicz [16], Becker, Bizri and Bresler [1] and Wickström [14] and [15]



Temperature distributions illustrating two dimensional heat flow in a prestressed concrete beam exposed to a standard fire calculated by stepwise solution of difference expressions. Weiss [13]. 300 - 600 °C $^{\circ}$ C

600 -

SOME SIMPLE EXACT SOLUTIONS



Example of an integral surface.

An exact solution to the one-dimensional case obeys the equation

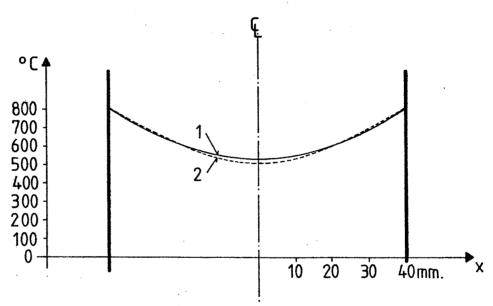
$$\frac{\partial \mathbf{T}}{\partial t} = \mathbf{a} \frac{\partial^2 \mathbf{T}}{\partial \mathbf{x}^2}$$

where a is the thermal diffusivity

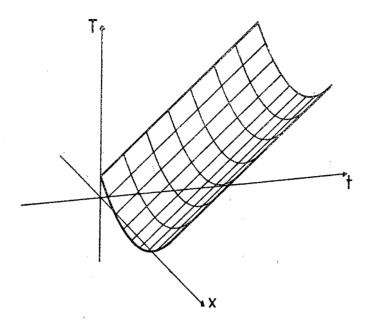
$$a = \frac{K}{\rho c_p}$$

In a x-t-T co-ordinate system such a solution will form a surface.

The inclination of the surface in the t-direction is at every point proportional to the derivative of the second order of the height T i.e. the approximate curvature of the x-direction.



Temperature distribution in the web of a prestressed beam after 1 h standard fire. 1) Parabola. 2) EDP calculation from Weiss [13].



Parabolic solution.

A surface of this kind happens to be an adequate conceptual tool for handling thermal problems related to fire exposed cross-sections.

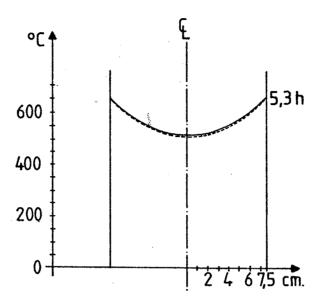
If the surface fulfils the boundary conditions, which for example consist of three boundary curves, it represents the particular integral that forms the exact solution to the problem.

If the temperature at the boundaries raises at constant speed 2aC₁ for example at two surfaces of a wall exposed from both sides, the parabola of the second degree is an exact solution

$$T = T_o + C_o x + C_1 x^2 + 2aC_1 t$$

where T_0 , C_0 and C_1 are arbitrary constants.

For a thin wall and a rapid heating this solution



is often applicable with a sufficient accuracy.

For cylindrical cross-sections the equation of conduction in polar co-ordinates becomes

$$\frac{1}{r} \frac{\partial T}{\partial r} + \frac{\partial^2 T}{\partial r^2} = \frac{1}{a} \frac{\partial T}{\partial t}$$

where r is the radius.

Also in this case the parabola of the second degree represents an exact solution

$$T = T_o + Cr^2 + 4aCt$$

Likewise does the parabola of the fourth degree

$$T = T_0 + C_1 r^2 + 4C_1 at + C_2 r^2 t + \frac{1}{16a} C_2 r^4 + 2C_2 at^2$$

An exponential solution also exist to the simple plane one-dimensional problem

$$T = T_o + C_o x + C_1 e^{(A_1 t - \sqrt{\frac{A_1}{a}}x)}$$

where T_0 , C_0 , C_1 and A_1 are arbitrary constants.

Furthermore the damped oscillation known from electromagnetism is an exact solution.

$$T = T_o + C_o x + C_1 e^{-\sqrt{v/2a}x} \sin(vt - \sqrt{v/2ax})$$

where $\mathbf{T}_{_{\mathrm{O}}}$, $\mathbf{C}_{_{\mathrm{O}}}$ and $\mathbf{C}_{_{1}}$ are arbitrary constants and \mathbf{v} is the angular velocity.

The surface temperature in this case must vary according to the expression

$$T = T_0 + C_1 \sin(vt)$$

If the surface temperature at time t=0 raises to a constant value T_{o} , a good approximate solution is known as

$$T = T_0 \left(1 - \frac{x}{3.363 \sqrt{at}} \right)^2$$

Finally a simple exact solution proposed by Joseph Fourier himself in [3] has to be mentioned.

Although it is not incorporated in the procedure presented on the following pages, it may be valuable for rough calculations in fire technology.

In fact it is an exact solution for a rectangular prism with a surface temperature varying exponentially in time.

The exponential decrease in temperature is interesting because it is easily superimposed to describe the variation according to the standard fire curve until the decay period (which unfortunately seldomly is described in the national standards).

As an example the present Danish Standard fire $\operatorname{curv}_{\varepsilon}$ is composed by exponential terms.

$$T - T_o = 1325-430e^{-0.2t}-270e^{-1.7t}-625e^{-19t}$$

With the same terminology as used before Fouriers solution for the surface temperature variation

 $T = C e^{-mt}$ in the three dimensional case i

$$T = C_0 e^{-mt} \cos(C_1 x) \cos(C_2 y) \cos(C_3 z)$$

where x, y and z are coordinates originating at the centre of the prism, and C_0 , C_1 , C_2 and C_3 are constants obeying the relation

$$m = a(C_1^2 + C_2^2 + C_3^2)$$

and demands concerning the values of the product of cosine functions at the surfaces of the prism.

Two- and one dimensional solutions are naturally for by introducing $C_3 = 0$ and $C_2 = 0$ respectively.

PRACTICAL APPROACH

The simple exact solutions mentioned in the previous chapter would be of limited interest if it was not for the fact that the equation of conduction allows superposition. That is, two exact solutions can be multiplied by constants and added, and the result is still an exact solution.

By means of this procedure many realistic boundary conditions can be fulfiled almost exactly.

In spite of the fact, that an uncertainty of about 50° C has to be accepted exclusively according to the difficulties in the assessment of the thermal conductivity, the author finds it convenient to use a simple procedure for the temperature calculation giving rise to uncertainties of about $30\text{-}40^{\circ}$ C.

The procedure proposed in this presentation approximates the surface temperature development of a fire exposed construction with an arbitrary rectangular cross-section to an idealized development that is achieved by superposition of solutions to the equation of conduction.

The surface temperature development is composed by three elementary developments representing three basic solutions. These are a fixed temperature superimposed by a harmonic oscillation and after a half period superimposed by an exponential solution in the cooling phase.

By means of these three simple elements every fire development can be simulated with a sufficient accuracy.

In this context the fire developments of interest are chiefly the standard fire defined by for example ISO 834, succeeded by realistic cooling sequences, or actual fire developments specified by the fire loads and the opening factors according to Magnusson and Thelandersson [9].

The simulated solution to the plane one-dimensional problem can be written as

$$T(x,t) = f_1(x,t) + f_2(x,t) + f_3(x,t)$$

where

$$f_1(x,t) = E'\left(1 - \frac{x}{3.363\sqrt{at}}\right)^2$$

$$\left(1 - \frac{x}{3.363\sqrt{at}}\right) > 0, \text{ else } f_1(x,t) = 0$$

and E' is the constant temperature.

$$f_2(x,t) = D'e^{-\sqrt{v/2ax}}\sin(vt-\sqrt{v/2ax})$$
for $(vt-\sqrt{v/2ax}) > 0$, else $f_2(x,t) = 0$

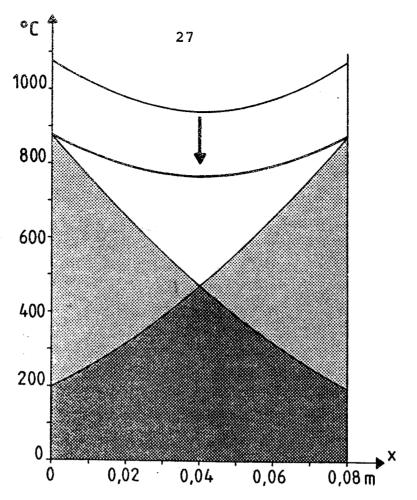
where D' is the amplitude of the harmonic oscillation at the surface, and with the half period called C', the angular velocity will be $v = \pi/C'$.

$$f_3(x,t) = \frac{D' + E'}{2(e^{LC'}-1)} \left(1 - e^{(L(t-C')-\sqrt{L/ax})}\right)$$

for
$$(L(t-C')-\sqrt{L/ax}) > 0$$
, else $f_3(x,t) = 0$

where
$$L = \frac{2}{C'} \ln \frac{3D'}{E'-2D'}$$
.

L defines the shape of the temperature curve in the cooling sequence being assessed by the surface temperatures:



Temperature distributions from two sides.

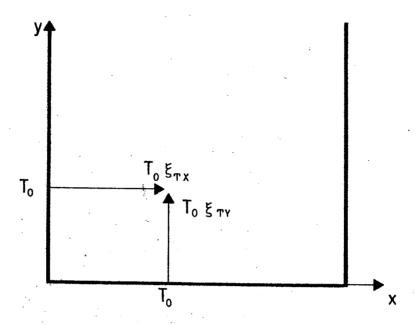
$$T(0,C') = E'$$

$$T(0,\frac{2}{3}C') = \frac{E'}{2}$$

$$T(0,2C') = \frac{E'-D'}{2}$$

For constructions simultaneously exposed to fire at two parallel surfaces, as for example walls and compression zones in top of beams, a one-dimensional solution from the one side is superimposed by the same solution from the other side, and the new temperature distribution is multiplied by the relation between the surface temperature wanted and the surface temperature from the added solution.

This simple procedure leads to temperature distributions which are in close accordance with known measured or precisely calculated distributions, as shown by the examples in Appendix A (calculated by Pedersen [11]).



Temperature reduction in a section exposed on three sides.

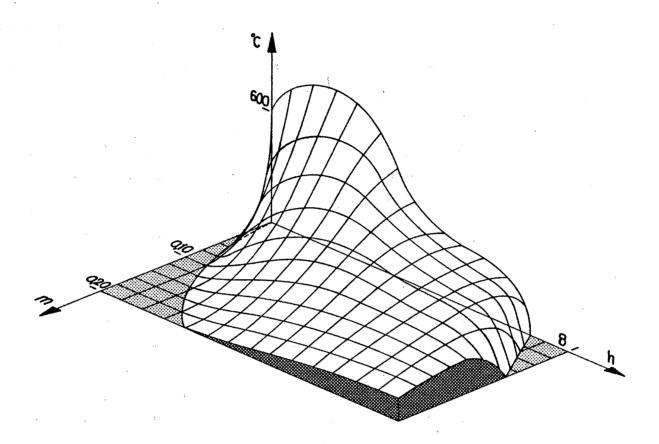
Analysing fire exposed concrete constructions it is often desirable being able to calculate the temperature in a point of a rectangular cross-section expos on the three sides.

Unlike the procedures dealt with on the previous pag a well known method is available for the simple calculation (for example Carslaw and Jaeger [2]).

If the reduction factor of the temperature is called ξ_{Tx} in the depth x of a section exposed at two sides and ξ_{Ty} in the depth y of a section exposed at one side, the temperature in the point (x,y) of the section exposed at three sides is calculated approximat as

$$T(x,y) = T_o(\xi_{Tx} + \xi_{Ty} - \xi_{tx}\xi_{Ty})$$

where T_{o} is the surface temperature.



Integral surface for the simple calculation procedure.

The author has developed a program for the simple calculation of temperatures in a cross-section exposed by fire at one-, two- or three surfaces by means of an advanced pocket calculator.

The program is listed in Appendix C, where also the documentation necessary for operating the program can be found.

Examples of temperature calculations of cross-sectic exposed on three sides are shown in Appendix B (calculated by Pedersen [11]).

It will be seen, that the simple calculation in this case leads to temperatures which are somewhat too high especially in the vicinity of a corner.

The increment of the temperatures in these zones of the cross-section is advantageous because it is of the same amount as the increment caused by the bevelled edges, that often are found on fire exposed rectangular construction elements as a result of the spalling effect.

This effect is caused by the flow of steam from the cross-section giving rise to an explosive destruction of the surface especially at convex corners where the thermal stresses are contributory to the phenomenon.

Because it happens at an early stage of the fire it causes certain changes of the isotherms of the cross section during the largest part of the fire development.

APPLICATIONS FOR THE SIMPLE CALCULATION PROCEDURE

It is obvious that a quick estimation of the temperature to a certain time at a point of a cross-section exposed to a certain fire is advantageous in relation to fire technological research.

The consulting engineer also has a need of such a procedure while selecting various sorts of fire pre-cautions or designing a concrete construction for fire resistance.

In this case especially the temperatures of the reinforcement bars are of interest, and the problem is characterized by the fact that they are located at single points of the cross-section.

Many of the fire technological phenomena of relevance for the designing process are far from simultaneous.

The time at which the maximum temperature occurs during a fire is highly dependent on the position of the point considered in the cross-section.

It is therefore important being able to maximize the temperature or temperature dependent phenomena at a reasonable speed and cost.

The procedure proposed is advantageous because it allows a mathematical treatment of many of these maximizations, or it simply offers a fast working subroutine for the temperature calculation as a part of a larger calculation.

The procedure is shown in Appendix D where it is translated into PL/1 (Programming Language One) for application in an EDP program.

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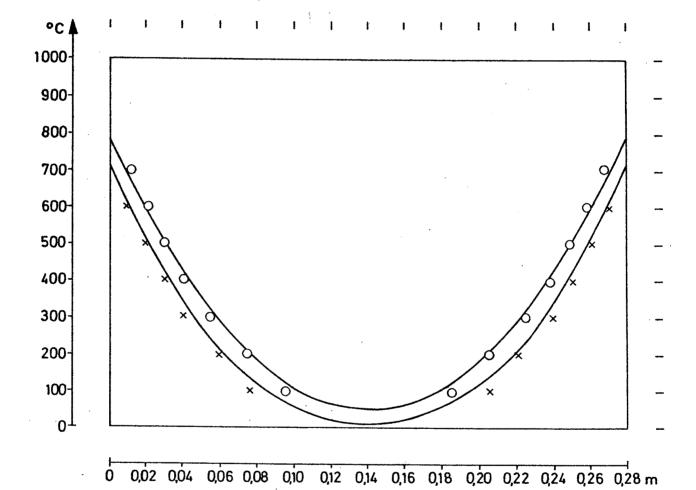
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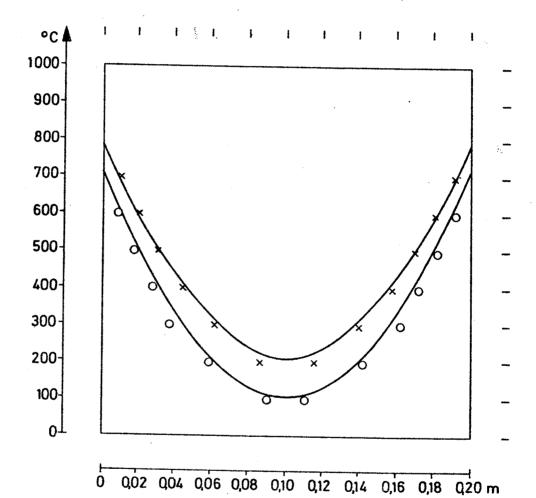
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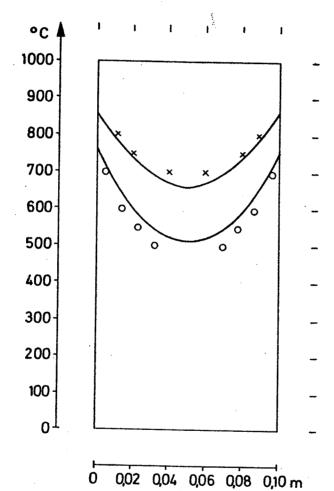
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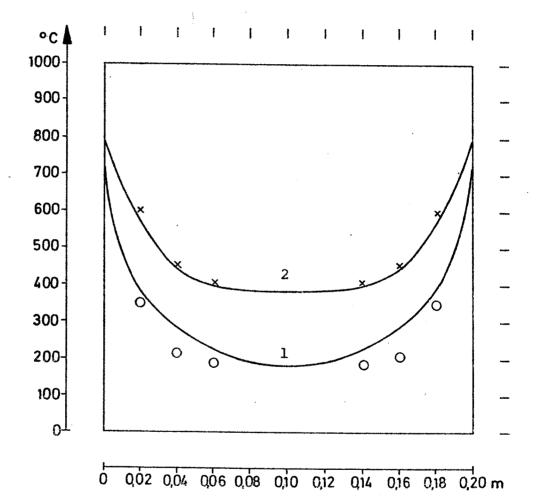
Temperature distribution calculated by the program Incendioret after 1.0 and 1.5 h standard fire. The points indicated are measured temperatures (Kordina et al. [6]).



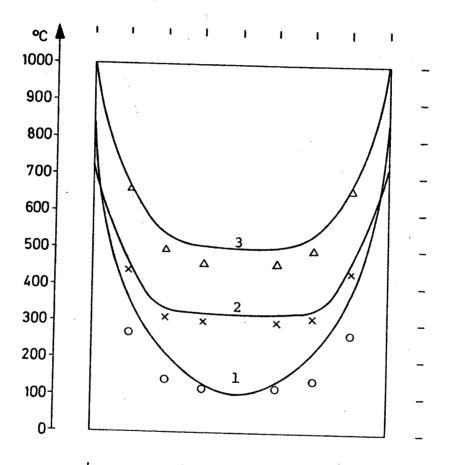
Temperature distribution calculated by the program Incendioret after 1.0 and 1.5 h standard fire. The points indicated are measured temperatures (Kordina et al. [6]).



Temperature distribution calculated by the program Incendioret after 1.0 and 1.5 h standard fire. The points indicated are measured temperatures (Kordina et al. [6]).



Distribution of maximal temperatures calculated by the program Incendioret at a fire of (opening factor, fire load) = 1: (0.06,150) (m²,MJ/m²) 2: (0.04,400) -The points indicated are temperatures calculated by EDP (Pettersson and Ödeen [12]) (TASEF-2 [15])



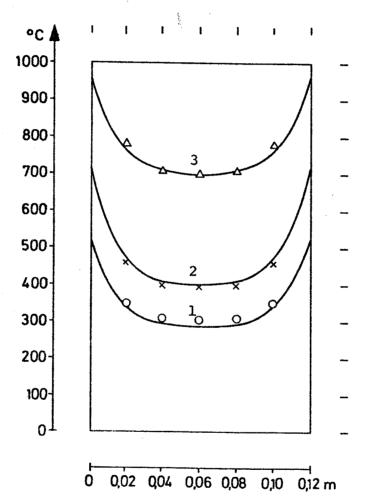
0,02 0,04 0,06 0,08 0,10 0,12 0,14 0,16 m

Distribution of maximal temperatures calculated by the program Incendioret at a fire of

(opening factor, fire load) = 1: (0.12,150) (m²,MJ/m²)
2: (0.04,200)

3: (0.12,900)

The points indicated are temperatures calculated by EDP (Pettersson and Ödeen [12]). (TASEF-2 [15])



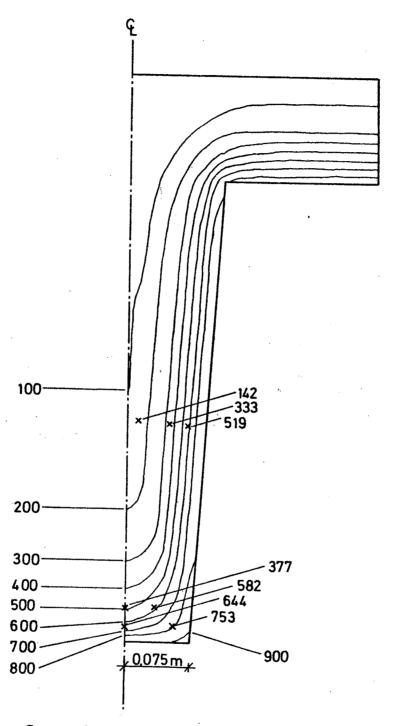
Distribution of maximal temperatures calculated by the program Incendioret at a fire of

(opening factor, fire load) = 1: (0.02,100) (m²,MJ/m²)
2: (0.04,200) - 3: (0.08,800) -

The points indicated are temperatures calculated by EDP (Pettersson and Ödeen [12])

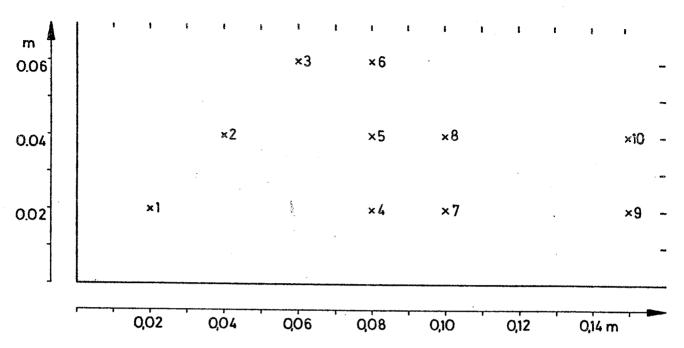
(TASEF-2 [15])

APPENDIX B



Comparison between temperatures calculated stepwise by means of difference expressions on EDP (Weiss [13]) - the isotherms - and by the program Incendioret - the points - in a beam exposed to 1 h standard fire.

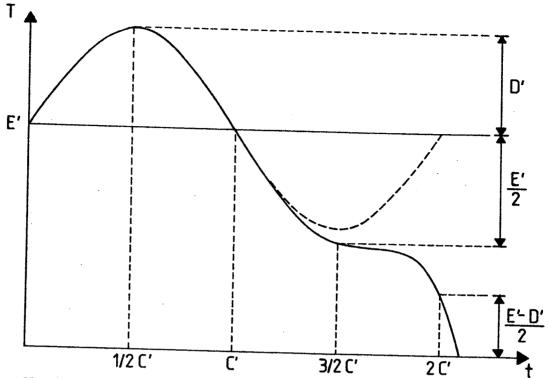
APPENDIX B



Comparison between maximal temperatures calculated by the program Incendioret and by EDP (Pettersson and Ödeen [12]) in a corner of a rectangular beam exposed to fires characterized by opening factor and fire load. (TASEF-2 [15])

(opening factor, fire load)	Beam with	Point	Temperature Incendioret	Temperature EDP
$(m^{\frac{1}{2}}, MJ/m^2)$	(m)		(°C)	(°c)
(0.04,200)	0.16	1 2 3 4 5	674 490 374 546 425	645 450 375 480 400
(0.04,300)	0.20	1 2 3 7 8	740 578 437 590 468	740 550 440 565 455
(0.06,600)	0.16	1 2 3 4 5 6	919 775 674 818 721 662	895 725 640 760 670 630
(0.08,600)	0.30	1 2 3 7 8 9	900 704 534 716 535 702 505	865 625 465 630 445 615 415





Variation of surface temperature for Incendioret.

THE POCKET CALCULATOR PROGRAM INCENDIORET

The pocket calculator program INCENDIORET is listed on the following pages. It is written for the pocket calculator TI Programmable 59 from Texas Instruments.

The program calculates the temperature in a point of a semiinfinite specimen, a slab exposed on two sides (and thus a slab with a perfect insolation on the one side calculated as a two side exposed slab of the double thickness) and a rectangular section exposed on the three sides.

The material is described by a fixed thermal diffusivity and the temperature variation on a surface exposed to the actual fire is described by three constants C', D' and E'.

The program is not able to calculate temperatures after the time of the maximum temperature at the point.

OPERATION OF THE PROGRAM INCENDIORET

First the material and the fire is described.

	TYPE		PRESS
	The thermal diffusivity	(m^2/s)	A¹
	The half period	(h)	C'
	The thermal amplitude	(°c)	D.
	The constant temperature	(°C)	E'
Then the	temperature can be calcula	ited in any	y point
	TYPE		PRESS
	The depth from the surface	e of	•
	a semiinfinite specimen	(m)	A
	The time in hour from the	start	
	of the fire	(h)	В
	s the temperature in a semecimen is a slab exposed o		
	TYPE.		PRESS
	The slab thickness	(m)	С
The displa	ay shows the new temperatu	ıre.	
	mperature in another depth		l at
the same t			
	TYPE		PRESS
	The new depth	(m)	D
If the sec	ction is rectangular then		
	TYPE		PRESS

The depth from the

(m)

E

third surface

INDATA FOR INCENDIORET

The values of D' and E' describe the temperature development at a plane surface absorbing radiant energy from the fire in a hemispherical space such as the surface of a slab or the bottom of a broad beam.

If the angle factor for radiation from the fire to the surface is less than 1.0, the values of D' and E' must be adjusted accordingly.

For web surfaces of beams at the ceiling of the enclosure it is recommended to multiply the values of D' and E' by the factor 0.9. These adjusted values are referred to as D' and E' web.

STANDARD FIRE

RATING	C¹	D'	E'	D' E' web
0.5 h	1.0	150	600	135 540
1.0 h	2.0	220	600	195 540
1.5 h	3.0	310	600	280 540
2.0 h	4.0	360	600	325 540
3.0 h	6.0	410	600	370 540 ·
4.0 h	8.0	460	600	410 540
OPENING FACTOR = $0.02 \text{ m}^{\frac{1}{2}}$				
FIRE LOAD	C'	D'	E '	D' E' web
75 MJ/m^2	0.6	100	360	90 330
100 MJ/m^2	1.3	185	390	170 350
150 MJ/m^2	2.2	250	400	225 360
200 MJ/m ²	3.2	310	400	280 360
250 MJ/m ²	3.6	350	400	315 360

OPENING FACTOR = $0.04 \text{ m}^{\frac{1}{2}}$

		•		•		
	FIRE LOAD	C†	D'	Ε°	D'	E web
•	75 MJ/m^2	0.4	135	470	120	420
•	100 MJ/m^2	0.7	200	510	180	460
	200 MJ/m^2	1.2	200	600	180	540
	300 MJ/m^2	1.8	240	600	215	540
		2.8	265	600	240	540
	500 MJ/m^2	3.4	300	600	270	540
OPENING FACTOR = 0	.06 m ¹ 2					
	FIRE LOAD	C i	D'	E'	D' web	E'web
	150 MJ/m^2	0.6	150	645	135	580
	300 MJ/m^2	1.2	195	680	1 7 5	610
	450 MJ/m^2	1.8	265	710	240	640
		2.4	270	745	245	670
	750 MJ/m ²	3.2	290	750	260	675
OPENING FACTOR = 0	.08 m ^½					
	FIRE LOAD	C'	D,	E '	D,	E'web
	100 MJ/m^2	0.3	145	645	130	580
	200 MJ/m^2	0.6	170	700	155	630
	400 MJ/m ²	1.2	200	725	180	650
	600 MJ/m^2	1.8	265	765	240	690
	800 MJ/m ²	2.4	270	800	245	720
	1000 MJ/m ²	3.0	290	805	260	725
OPENING FACTOR = 0	.12 m ^{.1} 2					
	FIRE LOAD	C'	D f	E'	D° web	E web
	150 MJ/m ²	0.3	220	710	200	640
	300 MJ/m^2	0.5	220	780	200	700
	600 MJ/m ²	1.0	250	800	225	720
	900 MJ/m^2	1.6	250	835	225	750
	1200 MJ/m ²	2.2	255	845	230	760
	1500 MJ/m ²	3.0	290	860	260	775

The pocket calculator program Incendioret

001234567890 0012344567890 0012344567890 001234567890 0055567890
76168L 108
0623445678901234456678900000000000000000000000000000000000
4331/X L9 4331/X L9 4331/X L9 53488 1 = 10 6353653 R 1 = 10 6
12234567890123345678901234456789012334567890123456789001234567890012345678900123456789001234567890012345678900123456789001234567890012345678900123456789001234567890012345678900123456789001234567890012345678900123456789000000000000000000000000000000000000
XL22 VYX L9

194 94 77 195 22 INV 196 23 LNX 197 94 +/- 198 85 + 199 01 1 200 95 = 201 65 × 202 43 RCL 203 25 25 204 95 = 205 42 STD 206 20 20 207 76 LBL 208 66 PRU 209 43 RCL 210 20 20 211 85 +	241 54) 242 85 + 243 01 1 244 54) 245 77 GE 247 55 ÷ 248 33 × 249 65 RCL 251 10 252 43 RCL 251 253 42 STD 253 42 STD 253 42 STD 253 42 STD 254 43 RCL 255 44 SUM 27 64 SUM 27 64 SUM 27 28 43 RCL 257 76 LBM 267 76 LBM 268 95 = 270 30 30 STD 268 269 42 STD 268 95 STD 271 30 STD 272 273 71 SBR 273 71 SBR 274 91 LBL 277 76 13 C 277 42 STD 278 43 RCL 277 278 43 RCL 277 279 43 RCL 277 279 43 RCL 288 41 1 1 RCL 287 288 41 1 1 RCL 288 42 STD 288 43 RCL 289 42 STD 289 43 RCL 289 44 RCL 280 44 RCL	301 43 RCL 302 18 18 303 65 2 304 022 2 305 95 23 306 42 STO 307 61 RCL 307 76 LBL 308 77 97 RCL 309 77 SBI 309 77 SBI 300 310 SBI 300 310 SBI 300
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361 36 RCL2 363 324+/- 364 853 RCL2 364 853 RCL2 365 843 RCL2 366 437 85 RCL2 367 911 R 12 SBR 370 111 RCL2 371 42 SBR 371 42 SBR 371 42 STD 373 374 42 STD 373 374 42 STD 377 85 RCL2 378 95 X RCL2 378 95 X RCL2 378 95 X RCL2 378 95 RCL2 379 85 RCL2 370 42 STD 371 SBR 371 SBR 372 42 STD 373 374 42 STD 373 374 42 STD 377 85 RCL2 378 43 RCL2 378 43 RCL2 379 42 RCL2 379 42 RCL2 411 85 RCL2 411 71 SBR 412 42 RCL2 411 71 SBR 413 42 RCL2 414 43 45 RCL2 415 427 428 RCL2 417 428 438 449 449 449 449 449 449 449 449 449 44		122345678901233456789012344567890123 442244224444444444444444444444444444	36 = 0 95 = 10 38 95 STD 38 91 R SS 91 R SS 91 R SS 91 R SS 91 R SS 91 R SS 95 R SS 95 R SS 95 R SS 95 R SS 96 SS 97 R SS 98 R SS 99 R SS 90 R SS 90 R SS 90 R SS 90 R SS 90 R SS 91 R SS 91 R SS 92 R SS 93 R SS 94 R SS 95 R SS 96 R SS 97 R SS 98 R
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APPENDIX D

PL/l Subroutine

```
TEMP: PROCEDURE(A3,C3,D3,E3,X3,T13);
     V = 3.141592654/C3;
      ARG2 = E3-2*D3;
      IF ARG2 < 0.02 THEN ARG2 = 0.02;
      ARG = 3*D3/ARG2;
      Z3 = 2*LOG(ARG)/C3;
      ARG1 = SQRT(V/(7200 \pm A3)) * X3;
      IF ARG1 = 0 THEN EX = 1;
      ELSE EX = EXP(-ARGI);
      ARG = V*TI3-ARG1;
      IF ARG > 0 THEN
      FUI = D3*EX*SIN(ARG);
      ELSE FU1 = 0;
      ARG = Z3*(T13-C3)-SQRT(Z3/(3600*A3))*X3;
      IF ARG > 0 THEN
      FU2 = (D3+E3)*(1-EXP(ARG))/(2*(EXP(Z3*C3)-1));
      ELSE FU2 = 0;
      ARG = 1-X3/(3.363*SQRT(A3*3600*TI3));
      IF ARG > 0 THEN
      FU3 = E3*ARG**2+FU2;
      ELSE FU3 = 0;
      RETURN (FU1+FU3);
      END TEMP:
```